

A Bistable Poly[2]catenane Forms Nanosuperstructures**

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Dedicated to David Reinhoudt on the occasion of his 65th birthday

The introduction of the mechanical bond^[1] into chemical compounds has contributed significantly to the realization^[2] of precise molecular architectonic tuning capabilities previously unattainable at the molecular scale. The relative molecular motions (e.g., circumrotation, translation, etc.) displayed within mechanically interlocked compounds, such as catenanes and rotaxanes, has led, as a result of templation^[3] during their synthesis, to the emergence of bistability^[4] necessary for the operation of molecular switches and machines^[2] with potential uses in molecular electronic devices (MEDs), and in nanoelectromechanical systems (NEMS) such as “mechanized” nanoparticles (MNPs)^[5] for controlled drug delivery. In the context of MEDs alone, molecular switches in the form of bistable rotaxanes have been used^[6] to pattern a 160-kbit molecular electronic memory at a density of 10^{11} bits per square centimeter in a two-terminal crossbar device that is smaller than the cross-section of a white blood cell.

Despite these achievements, the challenge remains to make these bistable, mechanically interlocked molecules not only readily accessible synthetically, but also to display materials processibility suitable for efficient device fabrication. Motivated by the emerging trends in materials chemistry, these highly programmable, mechanically interlocked compounds are now being incorporated into reticular networks^[7] and polymeric scaffolds.^[8] Mechanically interlocked polymeric scaffolds (MIPS), such as main-chain,^[10] pendent,

ant,^[11] and bridged^[12] classes of polycatenanes, however, remain a formidable synthetic challenge. MIPS have long been known^[7–12] to possess unique mechanical and dynamic properties. As a functional category of MIPS, mechanically interlocked switchable polymeric scaffolds^[13] (MISPS) could benefit enormously from the combination of bistable molecular switchability and materials processibility.

Recent research^[14] has uncovered efficient routes to synthesize the tetracationic cyclophane, namely CBPQT⁴⁺, or cyclobis(paraquat-*p*-phenylene)-containing compounds, using the Cu^I-catalyzed Huisgen 1,3-dipolar cycloaddition^[15] as the so-called click reaction.^[16] These findings have enabled convenient and modular syntheses of the required complex molecules in high yields.^[17] By incorporating a pendant alkyne group onto the π -electron deficient CBPQT⁴⁺ ring, the click reaction can be used to graft this tetracationic cyclophane, and (bistable) [2]catenanes containing it, onto polymers carrying azide side-chains. Herein, we report the synthesis, actuation, ground-state equilibrium thermodynamics and switching kinetics of bistable poly[2]catenane-based nanoparticles (NPs). These NPs give rise to hierarchical, self-assembled superstructures whose formation can be rationalized by electrostatic and surface effects.

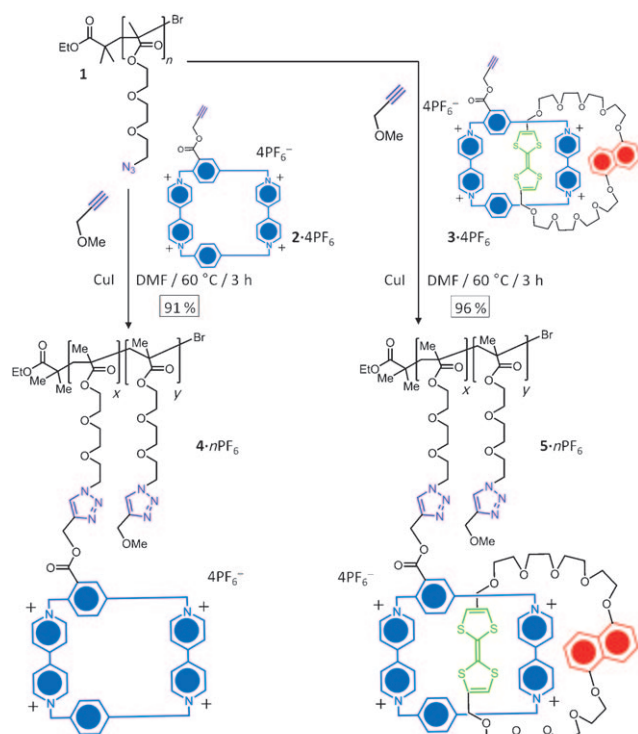
The copper-catalyzed click reaction between azide and alkyne functions in polymeric materials is reported^[18] to go to high conversion in part, it is believed, because of the increased solubility of the Cu^I species by the accumulating polytriazoles, which results in autocatalysis. The rates of click reactions involving alkynes added to azide-functionalized polymers have been shown^[19] to be much faster than those using the corresponding monomer, most likely as a result of the anchimeric assistance. Starting with the side-chain azide-functionalized polymer **1** ($M_w = 55\,000\text{ g mol}^{-1}$ and $M_n = 39\,000\text{ g mol}^{-1}$, PDI = 1.4), obtained by atom transfer radical polymerization (ATRP), 1) the bistable side-chain poly[2]-catenane co-polymer^[20] **5**·*n*PF₆, and the control co-polymer,^[20] and 2) the side-chain poly-CBPQT⁴⁺ **4**·*n*PF₆ were obtained (Scheme 1) in $\geq 95\%$ conversion, as determined by end-group analysis using ¹H NMR spectroscopy, and were isolated in 91 and 96 % yields, respectively. Assuming the click reaction proceeds with almost complete conversion and, based on the molar ratio of the azide functions to the alkyne-functionalized catenane **3**·4PF₆, a theoretical M_n of $128\,000\text{ g mol}^{-1}$ for **5**·*n*PF₆ was calculated; this value agrees with that estimated from end-group analysis by ¹H NMR spectroscopy. Analysis by size exclusion chromatography, coupled with detection by multi-angle light scattering (SEC-MALS), confirmed the

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[**] We gratefully acknowledge FENA and MURI for funding support. R.K. was supported by the MRSEC program of the National Science Foundation (DMR-0520513) at the Materials Research Center of Northwestern University. A.B.B. was supported by an NIH Postdoctoral Fellowship (1F32CA136148-01).

Supporting information for this article is available on the WWW under <http://dx.doi.org/10.1002/anie.200804558>.



Scheme 1. Click functionalization of polymer **1** ($n=160$) with the alkyne-functionalized tetracationic cyclophane **2·4PF₆** and the alkyne-functionalized bistable [2]catenane **3·4PF₆** for the formation of **4·nPF₆** ($x=27$, $y=133$) and **5·nPF₆** ($x=42$, $y=118$).

aggregation behavior of **5·nPF₆** in DMF with the calculated values for M_w of $(1.30 \pm 0.07) \times 10^6 \text{ g mol}^{-1}$, M_n of $(8.70 \pm 0.61) \times 10^5 \text{ g mol}^{-1}$, and a PDI of 1.5 ± 0.1 , corresponding to an RMS radius of gyration (r_g) of $127(\pm 2.5) \text{ nm}$ for the nanoparticles. The SEC-MALS conformation plot of $\log r_g$ as a function of $\log M_w$ for **5·nPF₆** yields a slope of $0.34 (\pm 0.01)$, a value which is consistent^[21] with spherical nanoparticles being present in solution (see Supporting Information for further details). Although SEC-MALS also reveals that **4·nPF₆** aggregates strongly in solution, the conformation plots do not indicate the presence of any spherical nanoparticles in solution. In a series of independent measurements,^[22] diffusion coefficients (D_0) of $(1.15 \pm 0.05) \times 10^{-12} \text{ m}^2 \text{ s}^{-1}$ for **5·nPF₆** and $(1.51 \pm 0.06) \times 10^{-9} \text{ m}^2 \text{ s}^{-1}$ for **3·4PF₆** were measured by double potential step chronocoulometry^[23] (see Supporting Information) in DMF at 298 K and applying the Cottrell relationship.^[24] Dynamic light scattering (DLS) measured a diffusion coefficient of $(2.22 \pm 0.21) \times 10^{-12} \text{ m}^2 \text{ s}^{-1}$ for **5·nPF₆**, and a hydrodynamic radius (R_h) of $174(\pm 32) \text{ nm}$, yielding an R_g/R_h value of $0.730(\pm 0.13)$ —again indicating^[25] the presence of spherical particles in solution. The [2]catenane and CBPQT⁴⁺ side-chain functionalized polymers were further characterized (see Supporting Information) by ¹H NMR and ¹H-¹H g-DQF-COSY NMR spectroscopies.

Bistable [2]catenanes, based on the highly preferred encirclement of the tetrathiafulvalene (TTF) over dioxynaphthalene (DNP) units, both located in the macrocyclic polyethers, by the CBPQT⁴⁺ ring have been investigated^[26] extensively. Switching within these molecules can be initi-

ated^[26,27] by either chemical or electrochemical stimuli. This switching manifests itself in the form of mechanical circumrotation of the macrocyclic polyether component with respect to the CBPQT⁴⁺ ring. The bistability presents itself as an equilibrium between two translational isomers in which the ground-state co-conformation (GSCC) and the metastable-state co-conformation (MSCC) correspond to the encirclement within the macrocyclic polyethers of the TTF and DNP units, respectively, by the CBPQT⁴⁺ ring. The ratio of GSCC to MSCC at 298 K is roughly 9:1, corresponding to a ΔG° value of $1.6 \text{ kcal mol}^{-1}$ when the CBPQT⁴⁺ ring is free to circumrotate between the TTF and DNP units in bistable catenanes in solution.^[28]

In solution, the thermochromic properties of both the **3·4PF₆**- and **5·nPF₆**-derived NPs were examined (Figure 1) by

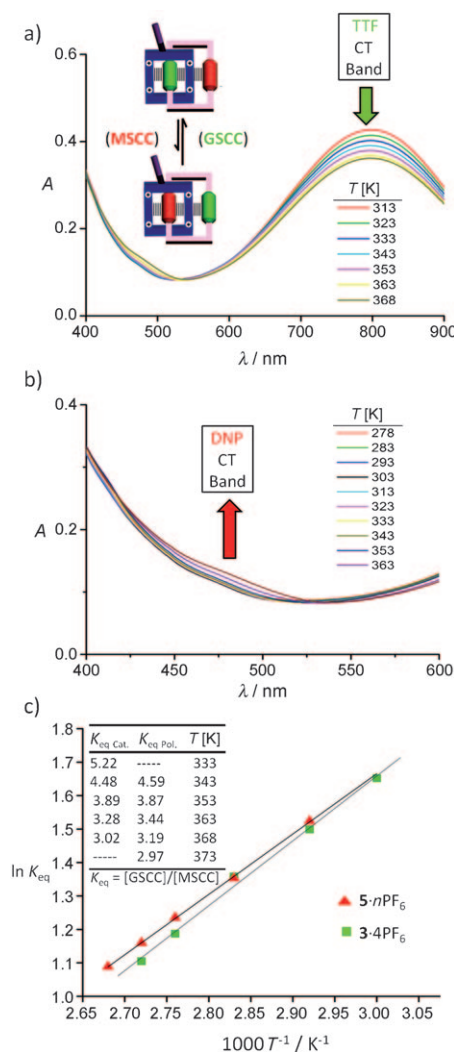


Figure 1. a) Variable temperature UV/Vis spectra of the alkyne-functionalized [2]catenane **3·4PF₆** and b) the side-chain poly[2]catenane **5·nPF₆** ($2 \times 10^{-4} \text{ M}$, DMF) illustrating the temperature-dependent equilibrium (inset) between the GSCC (TTF-CBPQT⁴⁺ charge-transfer band at 800 nm) and MSCC (DNP-CBPQT⁴⁺ charge-transfer band at 480 nm). c) Linear Van't Hoff plots for **3·4PF₆** ($R^2=0.996$) and **5·nPF₆** ($R^2=0.996$) for the thermally induced GSCC→MSCC equilibrium (inset); the ΔG°_{298K} , ΔH° , and ΔS° values for the process were obtained as described in the footnotes to Table 1.

variable temperature (VT) UV/Vis spectroscopy in DMF in order to quantify the ground-state equilibrium thermodynamic parameters. The charge-transfer (CT) band situated at 812 nm is characteristic^[26] of TTF residing inside CBPQT⁴⁺ (GSCC), whereas DNP encircled by CBPQT⁴⁺ (MSCC) produces a characteristic^[26] CT band at 475 nm in DMF. Upon cooling the solution, the TTF/CBPQT⁴⁺ CT band ($\epsilon = 4000 \text{ L mol}^{-1} \text{ cm}^{-1}$) increases in intensity while the DNP/CBPQT⁴⁺ CT band decreases, indicating that the equilibrium between the translational isomers favors the GSCC at lower temperatures. Heating the solution produces the opposite trend, namely an increase in the intensity of the DNP/CBPQT⁴⁺ CT band, accompanied by a decrease in the TTF/CBPQT⁴⁺ CT band. In conjunction with VT ¹H NMR spectra recorded to determine the equilibrium constant for the GSCC/MSCC ratio in **3-4PF₆** at 298 K, Van't Hoff plots were constructed in order to obtain the standard enthalpy (ΔH°) and standard entropy (ΔS°) for the equilibration (Table 1) of the GSCC with the MSCC. The GSCC→MSCC processes^[29] in **3-4PF₆** and **5-nPF₆** correspond to ΔG° values of $1.5(\pm 0.24)$ and $1.6(\pm 0.05) \text{ kcal mol}^{-1}$, respectively.

Tetrathiafulvalene undergoes a sequential and reversible two-electron oxidation process ($\text{TTF} \rightarrow \text{TTF}^+ \rightarrow \text{TTF}^{2+}$), generating cationic species, a property which triggers switching through circumrotation in **3-nPF₆** and **5-nPF₆** of the macrocyclic polyether with respect to the CBPQT⁴⁺ ring. Reduction of these cationic species back to being neutral forms the MSCC which re-equilibrates to the equilibrium mixture of the GSCC and MSCC. Switching of the GSCC to the MSCC, which was achieved electrochemically,^[30] was followed by cyclic voltammetry (CV), differential pulse voltammetry (DPV), and spectroelectrochemistry. In the case of the bistable poly[2]catenane **5-nPF₆**, the GSCC CT band at 812 nm bleaches (Figure 2a) as the voltage is increased and new absorptions emerge at 445 and 595 nm for the TTF^{•+} radical cation. At voltages above 700 mV, the simultaneous disappearance of the peaks at 445 and 595 nm and the appearance of a DNP CT band at 520 nm indicates the oxidation of TTF^{•+} to TTF²⁺ dication. Similar spectral changes are observed (Figure 2b) by spectroelectrochemistry when **5-nPF₆** is treated with 2 equivalents of $\text{Fe}(\text{ClO}_4)_3$. The TTF²⁺ dication, which is formed by this chemical oxidation, can be reduced back to a neutral TTF unit reversibly upon addition of aqueous ascorbic acid. These observations confirm that, even when the bistable [2]catenanes are attached to a polymer backbone and the poly[2]catenane assemblies to give larger secondary superstructures, they can still operate as

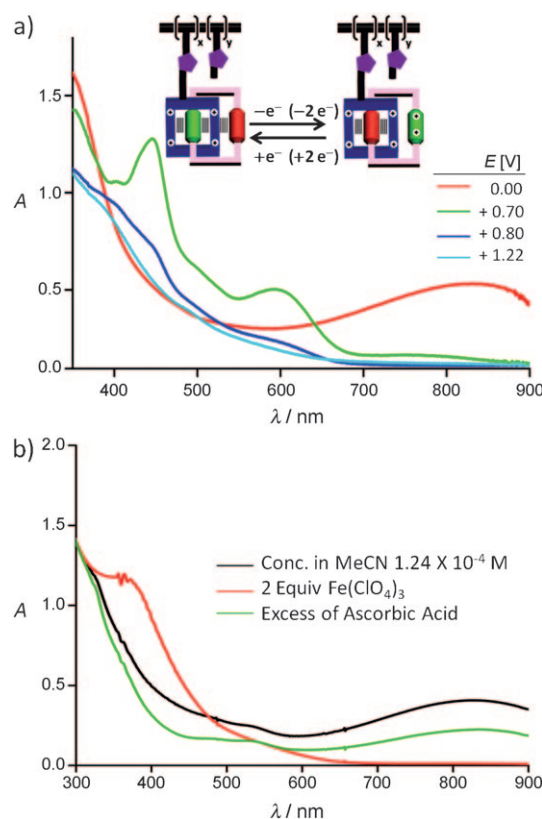


Figure 2. a) UV/Vis spectroelectrochemistry of the bistable poly[2]catenane **5-nPF₆** ($5 \times 10^{-4} \text{ M}$, MeCN, 0.1 M TBAPF₆, 298 K) recorded over a range of potentials illustrating the reversible two-electron electrochemically induced switching behavior of the bistable [2]catenane polymer side chains (inset). b) UV/Vis absorption spectra of the bistable poly[2]catenane **5-nPF₆** ($1.24 \times 10^{-4} \text{ M}$, MeCN, 298 K) illustrating the reversible chemically induced switching after the addition of 2 equivalents of the chemical oxidant $\text{Fe}(\text{ClO}_4)_3$ and reduction with an excess of ascorbic acid.

molecular switches under redox control (see the Supporting Information for CV measurements on **3-4PF₆** and **5-nPF₆**). Following electromechanical actuation, recovery of the MSCC/GSCC equilibrium is an activated process. The kinetics of the relaxation of the MSCC back to the GSCC can be quantified by variable scan rate CV (Figure 3a) and fitting (Figure 3b) a first order decay model to the population ratios of the metastable state and the relaxation times. The role of the physical environment on molecular electromechanical switching in solution,^[26,28,31] polymer gels,^[29] self-assembled monolayers,^[32] molecular switch tunnel junctions^[27,33] (MSTJs), and MEDs has already been probed^[28,34] in considerable detail. For **3-4PF₆** and **5-nPF₆**, the ΔG° values were found to be $17.8(\pm 5.0)$ and $17.4(\pm 0.5) \text{ kcal mol}^{-1}$, respectively, in DMF. We conclude that the bistable [2]catenane side-chains in **5-nPF₆** are equally accessible to solvent as in **3-4PF₆**, so both [2]catenanes experience very simi-

Table 1: Ground-state equilibrium thermodynamic parameters (ΔH° , ΔS° , $\Delta G^\circ_{298\text{K}}$) and switching kinetics ($\tau_{298\text{K}}$, $k_{298\text{K}}$, $\Delta G^\ddagger$) obtained after analysis of variable-temperature UV/Vis and variable-scan rate cyclic voltammetry experiments.

Species	ΔH° ^[a] [kcal mol ⁻¹]	ΔS° ^[b] [cal mol ⁻¹ K ⁻¹]	$\Delta G^\circ_{298\text{K}}$ ^[c] [kcal mol ⁻¹]	$\tau_{298\text{K}}$ ^[d] [s]	$k_{298\text{K}}$ ^[d] [s ⁻¹]	$\Delta G^\ddagger$ ^[d] [kcal mol ⁻¹]
3-4PF₆	$3.9(\pm 0.1)$	$7.9(\pm 0.2)$	$1.5(\pm 0.2)$	$1.9(\pm 0.9)$	$0.52(\pm 0.24)$	$17.8(\pm 5.0)$
5-nPF₆ ^[e]	$3.8(\pm 0.1)$	$7.5(\pm 0.1)$	$1.6(\pm 0.1)$	$0.9(\pm 0.1)$	$1.11(\pm 0.07)$	$17.4(\pm 0.5)$

[a] Obtained from a linear Van't Hoff plot where the slope of the line is equal to $-\Delta H^\circ/R$. [b] Obtained from a linear Van't Hoff plot where the intercept is equal to $\Delta S^\circ/R$. [c] Calculated from the Gibbs free energy $\Delta G^\circ = \Delta H^\circ - T\Delta S^\circ$. [d] Kinetic data obtained from 3 mm samples in DMF (0.1 M TBAPF₆), with a glassy carbon working electrode, at various scan rates versus Ag/AgCl.

lar environments. Perhaps the appended [2]catenanes in $5\text{-}n\text{PF}_6$ reside on the peripheries of the secondary superstructures in solution.

The self-assembled superstructures of $5\text{-}n\text{PF}_6$ were examined (Figure 4a–d) by scanning electron microscopy (SEM). In a highly dilute solution of $5\text{-}n\text{PF}_6$ in MeCN the polymers assemble into hollow spherical shells (see inset in Figure 4b). The outer radii of these structures are around 200 nm that is, in agreement with the results from SEC-MALS, DLS, and chronocoulometry experiments. In the dry state, the particles appear flattened/collapsed with holes present on their surfaces. The formation of these nanoscopic superstructures (NS) can be rationalized based on the minimization of electrostatic^[35] and surface energies^[36] during the assembly of multiple (n) $5\text{-}n\text{PF}_6$ charged “globules” into one NS.

For both of these structures, the electrostatic energy can be expressed as

$$E_{\phi,i} = \frac{1}{2} \int_V \rho_+ \Phi \, dV$$

where ρ_+ is the charge density inside of the polymer, and the potential, Φ , is estimated using the Poisson equation,

$$\nabla^2 \Phi = \frac{-\rho_+}{\epsilon_0 \epsilon_p},$$

for the polymer phase and the linearized form of the Poisson–Boltzmann equation,

$$\nabla^2 \Phi = \kappa^2 \Phi$$

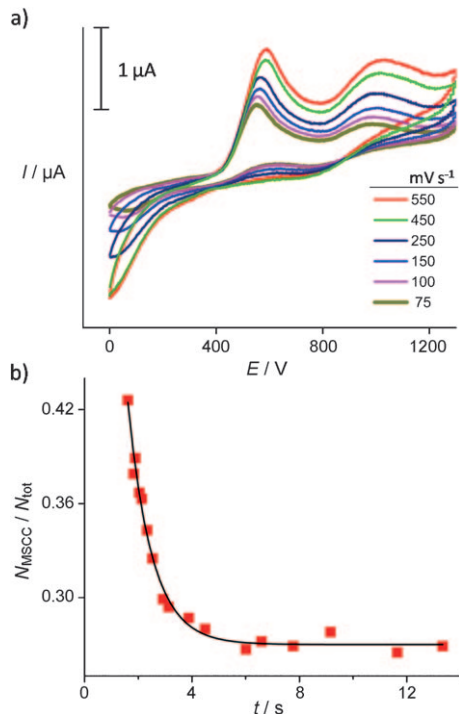


Figure 3. a) Series of second cycle CVs for $5\text{-}n\text{PF}_6$ taken at varying scan rates ($75/550 \text{ mV s}^{-1}$; $1 \text{ M TBAPF}_6/298 \text{ K/DMF}$ /versus Ag/AgCl). b) Fitting of a first-order decay model to the population ratios of the metastable state and relaxation times.

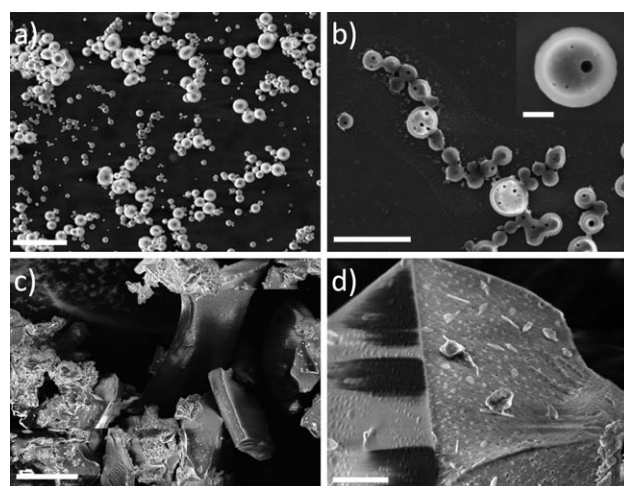


Figure 4. a,b) SEM images of hollow polymer nanosized superstructures prepared by drop-casting a dilute solution of $5\text{-}n\text{PF}_6$, followed by removal of solvent under vacuum. SEM image (b, inset) of a single hollow polymer superstructure with clearly visible pores in the shell. c,d) SEM images of larger assembled superstructures prepared by slow solvent evaporation of $5\text{-}n\text{PF}_6$ on glass. The scale bars are a) $4 \mu\text{m}$, b) $2 \mu\text{m}$ and 200 nm (inset), c) $20 \mu\text{m}$, and d) $2 \mu\text{m}$.

for the surrounding solution (κ^{-1} is the Debye screening length, ϵ_0 is the permittivity of free space, and ϵ_p is the dielectric constant of the NS). The surface energy is calculated from $E_{\sigma,i} = \sigma A$, where σ is the surface tension of the polymer–solution interface and A is the surface area. Calculations detailed in the Supporting Information show that 1) the process of self-organization is indeed energetically favorable and 2) the dimensions of the NS can be found by minimizing the energetic change of this process per globule,

$$\frac{\Delta E}{n} = \frac{(E_{\phi,\text{NS}} + E_{\sigma,\text{NS}})}{n} - E_{\phi,5\text{-}n\text{PF}_6} - E_{\sigma,5\text{-}n\text{PF}_6},$$

with respect to the NS radius. Specifically, for an experimentally observed shell thickness of 20 nm, and for physically reasonable values of other parameters ($\epsilon_p \approx 37$, $\sigma \approx 10 \text{ mN m}^{-1}$, $\kappa^{-1} \approx 1.6 \text{ nm}$, and $\rho_+ \approx 10^5 \text{ C m}^{-3}$), the model predicts an outer NS radius of 181 nm—that is, close to the experimental value. Finally, we observe that upon slow solvent evaporation on glass (Figure 4c,d), the polymeric NPs can form higher-level aggregates such as chains or joist-like chunks. Assembly of these structures is most likely driven by van der Waals forces that become increasingly important in the case of larger particles.

In summary, we have developed a highly efficient synthetic route to bistable side-chain poly[2]catenanes that self-assemble into novel types of supramolecular architectures including hollow spheres and their higher aggregates. These structures form as a result of a subtle interplay between electrostatic and surface forces acting as one in the system. A combination of analytical tools provides clear evidence that electromechanical switching of bistable [2]catenane units occurs in both free polymers and their nanoscopic assemblies. The maintenance of function in bistable mechanically interlocked units across ever-increasing levels of structural and

superstructural complexity argues well for the emergence of mechanically interlocked switchable polymeric scaffolds, i.e., MISPS for use in molecular electronic devices.

Received: September 16, 2008

Revised: November 4, 2008

Published online: January 29, 2009

Keywords: molecular devices · molecular switches · polymers · self-assembly · template synthesis

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